## Météo-France T-matrix code documentation

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November 13, 2020

## Contents

1	Introduction				
2	Program Tmatrix.f				
	2.1	Steps of the program			
	2.2	Option	ns and equations for aspect ratio, diameters, dieletric constant and oscillation		
		2.2.1	Aspect ratio r	3	
		2.2.2	Maximum diameter $D_m$	5	
		2.2.3	Description of the melting process and estimation of $D_{eq}$ and $D_{eqr}$	5	
		2.2.4	Wet hydrometeor content and mass water fraction	9	
		2.2.5	Dielectric constant EPSX	9	
		2.2.6	Compiling	12	
		2.2.7	Running the program	12	
3	Program TmatInt.f90				
	3.1 Steps of the program			13	
3.2 Compiling		Comp	iling	15	
	3.3	3 Running the program			
R	efere	nces		15	

## 1 Introduction

The T-matrix method is a computational technique of light scattering by nonspherical particles originally formulated by Waterman (1965). The T-Matrix ( $\mathbf{T}$  for Transition-matrix) allows to calculate the scattered field for a given incident field. More explanations are given for example in the PhD manuscrit section 1.2.1 of Le Bastard (2019)). The dual-polarization radar variables can then be calculated from the scattering coefficients, for a given radar wavelength.

The code used at Météo France is derived from Mishchenko and Travis (1994) code and has originally been downloaded from:

http://www.giss.nasa.gov/staff/mmishchenko/t\_matrix.html

The version used is : Extended-precision T-matrix code for nonspherical particles in a fixed orientation: amplq.lp.f, lpq.f, and amplq.par.f In order to run the code, it is necessary to install LAPACK packages:

http://www.netlib.org/lapack/

The Météo-France T-matrix version includes 2 programs:

- Tmatrix.f (fortran 77): computation of the scattering coefficients for a single hydrometeor with a given shape/diameter/dielectric constant
- TmatInt.f90 (fortran 90): integration of the scattering coefficients over the particle size distribution

This code was first created at Météo-France by Al-Sakka *et al.* (2013) and has been continually upgraded, and used for the studies of Augros *et al.* (2016), Taufour *et al.* (2018), Borderies *et al.* (2018), Le Bastard *et al.* (2019) and Thomas *et al.* (2020).

## 2 Program Tmatrix.f

This program computes back- and forward scattering coefficients (from which one can calculate the polarimetric variables) for a given radar wavelength. The coefficients are computed for different particle sizes. They are function of elevation angle, radar wavelength, temperature, liquid water fraction and hydrometeor type. The program parameters are defined in a parameter file, which is read in the beginning of the program. Calculation of scattering coefficients is made by the procedure TMD.

Out of this procedure, the scattering coefficients (S11, S22, S11f, S22f) and DPOL variables are written in two different files, for a range a temperature values, elevation angles, diameters and liquid water fraction (for graupel).

#### 2.1 Steps of the program

- 1. Reading of the hydrometeor constants file: contains  $a_j$  and  $b_j$  coefficients for the mass diameter relation:  $m(D) = a_j * D^{b_j}$
- 2. Reading of the parameter file "TmatParam BandType" (see example in Figure 1):

contains the hydrometeor type, the options for oscillation (oscill and sigbeta), for the dieletric constant (DIEL), for the aspect ratio (CEPS) and the min/step/max ranges of the varying parameters: radar wavelength (LAM), radar elevation (ELEV), temperature in degree C (Tc), exponent (expD) of the spherical equivalent diameter (D), and mass water fraction  $(F_w)$ .

3. Loops over LAM, Tc, ELEV,  $F_w$  and D (spherical equivalent diameter)

Calculation of:

- aspect ratio  $r = \frac{1}{EPS}$
- maximum diameter  $D_m$
- partially melted equivalent diameter  $(D_{eq} : \text{needed as input of the TMD subroutine})$ and fully melted equivalent diameter  $(D_{eqr})$
- dielectric constant EPSX
- Call of procedure TMD, which retrieves the scattering coefficients S11, S22, S11f, S22f as a function of LAM, ELEV, Tc, D<sub>eq</sub>, oscill, sigbeta, EPSX, EPS
- 5. Calculation of DPOL variables from the T-matrix scattering coefficients using the equations detailed in Appendix of Augros *et al.* (2016).

	Duvrir 🔻 主	TmatParam_Srr ~/Programmes/T		
1	type:rr			
2	oscill:0			
3	sigbeta:1			
4	DIEL:1			
5	CEPS:1			
6	LAMmin:106.2			
7	LAMmax:106.2			
8	LAMstep:0.1			
9	ELEVmin:0.0			
10	ELEVmax:20.0			
11	ELEVstep:4.0			
12	Tcmin:-20.0			
13	Tcmax:40.0			
14	Tcstep:1.0			
15	expDmin:0.0			
16	expDmax:3.0			
17	expDstep:0.1			
18	Fwmin:0.0			
19	Fwstep:0.1			
20	Fwmax:0.0			

Figure 1: Example of a parameter file for rain at S-band

- 6. Calculation of DPOL variables (ZhhR, ZvvR et KdpR) using the theory of Rayleigh for spheroids. The corresponding equations are detailed in section 3.4.3 of Caumont (2007).
- 7. Writing of the scattering coefficients and DPOL variables in two distinct files: TmatCoefDiff\_Srr and TmatVarPol\_Srr (example for rain at S band)

# 2.2 Options and equations for aspect ratio, diameters, dieletric constant and oscillation

### **2.2.1** Aspect ratio r

The option for aspect ratio r is determined by parameter CEPS.

The variable needed as input of the T-matrix TMD routine is:

$$EPS = \frac{1}{r}$$

1. Rain

CEPS = 1

The aspect ratio of raindrops  $r_r$  follows the formulation of Brandes *et al.* (2002), as in Ryzhkov *et al.* (2011).

$$r_r = 0.9951 + 0.02510D - 0.03644D^2 + 0.005303D^3 - 0.0002492D^4$$
 (1)

with D the spherical equivalent diameter in mm

#### 2. Primary ice

In a first approximation, pristine ice crystals are modelled as spheres because of their random orientation, as done in Caumont *et al.* (2006).

CEPS = 5

$$r_i = 1 \tag{2}$$

#### 3. **Snow**

CEPS = 6

• Dry snow

The aspect ratio of dry aggregated snow varies between 0.6 and 0.8 Straka (2000). In this code, it is assumed to be 0.75 as in Jung *et al.* (2008). But to avoid higher  $Z_{dr}$  for low diameters due to higher snow density, a linear decrease in the axis ratio from 1 to 0.75 for diameters from 0 to 8 mm and a constant axis ratio of 0.75 for diameters higher than 8 mm is simulated (see section 3.3.2 of Augros *et al.* (2016)):

$$r_s = \begin{cases} 1 - \frac{1 - 0.75}{8 - 0} D & \text{si } D \le 8 \text{ mm} \\ 0.75 & \text{si } D > 8 \text{ mm} \end{cases}$$
(3)

• Wet snow

In Augros *et al.* (2016), snow is considered dry only because in the ICE3 microphysics scheme, graupel is the only ice species that can have a wet growth mode. When the air temperature is warmer than 0 °C, small primary ice crystals are immediately converted into cloud water, snowflakes are transferred into the graupel category and finally graupel particles melt by shedding all the liquid water into raindrops (see Appendix of Lascaux *et al.* (2006)).

However, in the T-matrix code, the parametrization of wet snow is possible, if a dedicated parameter file for wet snow (TmatParam\_Sws) is created with a range of  $F_w$  values between 0 and 1.

In that case, the axis ratio of wet snow is parametrized as a combination of the axis ratios of dry snow and rain, following equation (14) of Ryzhkov *et al.* (2011):

$$r_{ws} = r_s + F_w(r_r - r_s) \tag{4}$$

In this equation, the axis ratio of rain  $r_r$  is estimated with Equation 1 but the spherical equivalent diameter D is replaced by the fully melted equivalent diameter  $D_{eqr}$  (see estimation in 2.2.3).

#### 4. Graupel

CEPS = 9

• Dry Graupel

The aspect ratio of dry graupel  $r_q$  is the one proposed by Ryzhkov *et al.* (2011):

$$r_g = \begin{cases} 1 - 0.02D & \text{if } D \le 10 \text{ mm} \\ 0.8 & \text{if } D > 10 \text{ mm} \end{cases}$$
(5)

 $r_g = 0.8$  mm for D > 10 mm in the study of Le Bastard (2019) but  $r_g = 0.85$  mm for D > 10 mm in the study of Augros *et al.* (2016) to account for oscillation that is neglected as explained in section 3.3.3 of Augros *et al.* (2016).

• Wet graupel

The aspect ratio of wet graupel  $r_{wg}$  follows the formulation of Ryzhkov *et al.* (2011), deduced from the experimental results of Rasmussen *et al.* (1984):

$$r_{wg} = \begin{cases} r_g - 5.0(r_g - 0.8)F_w & \text{if } F_w \le 0.2\\ 0.88 - 0.40F_w & \text{if } 0.2 < F_w \le 0.8\\ 2.8 - 4.0r_w + 5.0(r_w - 0.56)F_w & \text{if } F_w > 0.8 \end{cases}$$
(6)

where  $r_w$  is the aspect ratio of a raindrop of diameter  $D_{eqr}$ , which is produced as a result of the graupel melting.

#### **2.2.2** Maximum diameter $D_m$

The maximum diameter  $D_m$  of the spheroid is expressed as a function of the spherical equivalent diameter D using:

$$D_m = D \ EPS^{\frac{1}{3}}$$

$$D_m = D \ r^{-\frac{1}{3}}$$
(7)

#### 2.2.3 Description of the melting process and estimation of $D_{eq}$ and $D_{eqr}$

In MesoNH, graupel is the only specie that can have a wet growth mode or that can be melting (see Lascaux *et al.* (2006)). However, in the T-matrix code, the parametrization of wet snow is possible as proposed by Le Bastard (2019), if a dedicated parameter file for wet snow is created with a range of  $F_w$  values between 0 and 1 (see their chapter 4.1).

 $D_{eq}$  and  $D_{eqr}$  are the partially and fully melted equivalent diameters.  $D_{eq}$  is needed as input of the TMD subroutine to compute the scattering coefficients.  $D_{eqr}$  is needed for the melting species to estimate their axis ratio.  $D_{eqr}$  is also used in the integration of the scattering coefficients over the full Particule Size Distribution (PSD) that is done in TmatInt.f90 (see section 3). It is thus estimated for all species even those that are not potentially melted species (rain and primary ice).

#### 1. Rain

$$\boxed{D_{eq} = D_{eqr} = D} \tag{8}$$

#### 2. Primary ice

For primary ice, the liquid water fraction  $F_w$  is set to 0:

$$D_{eq} = D \tag{9}$$

The fully melted equivalent diameter  $D_{eqr}$  is calculated as a function of the density of liquid water  $\rho_w$  (1000 kg m<sup>-3</sup>) and primary ice  $\rho_i$  (kg m<sup>-3</sup>)

$$D_{eqr} = \left(\frac{\rho_i}{\rho_w}\right)^{\frac{1}{3}} D \tag{10}$$

with

$$\rho_i = \frac{a_i D_m^{b_i}}{\frac{\pi}{6} D^3} \tag{11}$$

 $a_i = 0.82$  and  $b_i = 2.5$  are the ICE3 mass-diameter coefficients for primary ice. Equation 10 is also used by Wolfensberger and Berne (2018).

#### 3. Snow

In Augros *et al.* (2016), the liquid water fraction  $F_w$  is set to 0:  $D_{eq} = D$ 

But in the T-matrix code, the parametrization of a wet snow specie is possible, following Le Bastard (2019). In that case, the partially melted equivalent diameter  $D_{eq}$  is calculated as:

$$D_{eq} = \left(\frac{\rho_s}{\rho_{ws}}\right)^{\frac{1}{3}} D \tag{12}$$

where the density of wet snow  $\rho_{ws}$  (kg m<sup>-3</sup>) is estimated as a function of the mass water fraction  $F_w$  following equation (4) of Jung *et al.* (2008):

$$\rho_{ws} = \rho_s (1 - F_w^2) + \rho_w F_w^2 \tag{13}$$

 $\rho_s$  is calculated like for primary ice with Equation 11 with  $a_s = 0.02$  and  $b_s = 1.9$ .

 $D_{eqr}$  is calculated like for primary ice with Equation 10 by replacing  $\rho_i$  by  $\rho_s$ .

#### 4. Graupel

Ryzhkov *et al.* (2011) and Rasmussen and Heymsfield (1987) suggest that if the initial density of the dry hail or graupel is less than the one of a solid ice, water first soaks particle interior and fills all air cavities. Once melted water saturates the particle interior  $(F_w = F_{w_{sat}})$ , it starts building a water shell.

The different steps of the melting process are described hereafter following Le Bastard (2019) and illustrated in Figure 2 (see also Figure 1 of Rasmussen and Heymsfield (1987)).

#### • Step 0: $F_w = 0$

The initial volume of air cavities  $V_c^{(i)}$  within the dry graupel ( $F_w = 0$ ) can be expressed using:

$$V_c^{(i)} = V_a - V_i$$

with  $V_i$  the volume occupied by pure ice

$$V_i = \frac{m_i}{\rho_i} = \frac{m_g}{\rho_i} = \frac{\pi}{6} D^3 \frac{\rho_g}{\rho_i} \tag{14}$$

with  $\rho_i$  the density of pure ice (916 kg m<sup>3</sup>)



Figure 2: Illustration of the graupel melting process as a function of mass water fraction  $F_w$ . Air (A), ice (I) and liquid water (W) are shown respectively in grey, white, and blue. The various combinations [matrix, inclusions] used to calculate the dielectric constant according to the Maxwell-Garnett approximation (see 2.2.5) are also shown. Taken from Figure 4.1 of Le Bastard (2019).

 $(m_g = m_c + m_i = m_i$  because the mass of air cavities  $m_c = 0)$  $V_c^{(i)}$  can thus be expressed as:

$$V_c^{(i)} = \frac{\pi}{6} D^3 (1 - \frac{\rho_g}{\rho_i}) \tag{15}$$

• Step 1:  $0 < F_w \leq F_{w_{sat}}$ 

When the graupel starts melting, the melted water first soaks the air cavities reducing the particle volume  $V_{wg}$ :

$$V_{wg} = \frac{\pi}{6} D_{eq}^{3}$$
  
=  $\frac{\pi}{6} D^{3} - \frac{F_{w} m_{g}}{\rho_{g}}$   
$$V_{wg} = \frac{\pi}{6} D^{3} (1 - F_{w})$$
 (16)

The equivalent diameter of the wet graupel  $D_{eq}$  is then:

$$D_{eq} = (1 - F_w)^{\frac{1}{3}}D$$
(17)

The volume of liquid water  $V_w$  formed due to the ice melting can be expressed as a function of the mass water fraction  $F_w$ , the particle mass  $m_{wg} = m_g$  and the density of liquid water  $\rho_w$ :

$$V_w = \frac{F_w m_g}{\rho_w}$$
$$V_w = \frac{\pi}{6} D^3 F_w \frac{\rho_g}{\rho_w}$$
(18)

As the graupel melts, the density of the cavities within the graupel is constant but their total volume (including both air and water) is reduced as is the graupel volume. The volume of the cavities can be expressed as:

$$V_c = \frac{\pi}{6} D_{eq}^3 (1 - \frac{\rho_g}{\rho_i})$$
(19)

The cavities are fully soaked when  $V_c = V_w$ .  $F_{w_{sat}}$  can thus be estimated using:

$$\frac{\pi}{6} D_{eq}^{3} (1 - \frac{\rho_{g}}{\rho_{i}}) = \frac{\pi}{6} D^{3} F_{w_{sat}} \frac{\rho_{g}}{\rho_{w}}$$

$$(1 - F_{w_{sat}}) D^{3} (1 - \frac{\rho_{g}}{\rho_{i}}) = D^{3} F_{w_{sat}} \frac{\rho_{g}}{\rho_{w}}$$

$$1 - \frac{\rho_{g}}{\rho_{i}} = F_{w_{sat}} (\frac{\rho_{g}}{\rho_{w}} + 1 - \frac{\rho_{g}}{\rho_{i}})$$

$$F_{w_{sat}} = \frac{1 - \frac{\rho_{g}}{\rho_{i}}}{1 - \frac{\rho_{g}}{\rho_{i}} + \frac{\rho_{g}}{\rho_{w}}}$$

$$\boxed{F_{w_{sat}} = \frac{\frac{1}{\rho_{i}} - \frac{1}{\rho_{g}}}{\frac{1}{\rho_{i}} - \frac{1}{\rho_{g}} - \frac{1}{\rho_{w}}}}$$
(20)

• Step 2:  $F_w > F_{w_{sat}}$ 

Once fully soaked, the particle starts building a water shell.

At this stage, the graupel is composed of pure ice and water only. Its mass  $m_{wg}$  can be expressed as a function of the mass of ice  $m_i$  + the mass of water  $m_{w_{tot}}$ , including the water within the core and in the shell:

$$m_{wg} = m_g = m_i + m_{w_{tot}} \tag{21}$$

Thus, its volume  $V_{wg}$  is:

$$V_{wg} = \frac{m_i}{\rho_i} + \frac{m_{w_{tot}}}{\rho_w}$$
$$= m_g \left(\frac{1 - F_w}{\rho_i} + \frac{F_w}{\rho_w}\right)$$
$$= \frac{\pi}{6} D^3 \rho_g \left(\frac{1 - F_w}{\rho_i} + \frac{F_w}{\rho_w}\right)$$

We can deduce the equivalent diameter  $D_{eq}$ :

$$D_{eq} = \left[\rho_g \left(\frac{1 - F_w}{\rho_i} + \frac{F_w}{\rho_w}\right)\right]^{\frac{1}{3}} D$$
(22)

For graupel, the fully melted equivalent diameter  $D_{eqr}$  is calculated like for primary ice and snow with Equation 10 by replacing  $\rho_i$  by  $\rho_g$ .

#### 2.2.4 Wet hydrometeor content and mass water fraction

As in Jung *et al.* (2008), graupel is assumed to be wet when graupel and rainwater coexist. Two options have been parametrized in the python forward operator program (the program that reads the output of the T-matrix code) to estimate the wet graupel content  $M_{wq}$  (kg m<sup>-3</sup>):

1. The rainwater content  $M_r$  (kg m<sup>-3</sup>) is added to the graupel content  $M_g$ , as done in Le Bastard (2019) and in Wolfensberger and Berne (2018):

$$M_{wg} = M_g + M_r \tag{23}$$

 $M_g$  and  $M_r$  are set to 0

2. Only the dry graupel content  $M_{wg}$  is transferred to the wet graupel content, as done in Augros *et al.* (2016) and Jung *et al.* (2008):

$$M_{wg} = M_g \tag{24}$$

 $M_q$  is set to 0 and rainwater coexists with wet graupel

In both cases, the mass water fraction within wet graupel  $F_w$  is estimated as:

$$F_w = \frac{M_r}{M_g + M_r} \tag{25}$$

#### 2.2.5 Dielectric constant EPSX

#### 1. **Rain**

DIEL=1

The Debye model is used for rain as in Caumont *et al.* (2006):

$$\epsilon_w = \epsilon_\infty + \frac{\epsilon_0 - \epsilon_\infty}{1 - if/f_D}, \qquad (26)$$

where  $\epsilon_0$  and  $\epsilon_{\infty}$  are respectively the dielectric static and high frequency coefficients and f is the radar frequency (Hz), and  $f_D$  is the relaxation frequency (Hz). Liebe *et al.* (1991) gave the following values for these parameters :

$$\epsilon_0 = 77,66 - 103,3\Theta \tag{27}$$

$$\epsilon_{\infty} = 0,066\epsilon_0 \tag{28}$$

$$f_D = (20, 27 + 146, 5\Theta + 314\Theta^2)10^9 \tag{29}$$

where the inverse relative temperature  $\Theta$  is related to T (K) with the following formula:

$$\Theta = 1 - \frac{300}{T} \tag{30}$$

#### 2. Snow, graupel and ice crystals

The effective permittivity  $\epsilon_{eff}$  of non homogeneous hydrometeors can be estimated based on the Maxwell-Garnett approximation (Maxwell Garnett, 1904), considering a matrix medium with permittivity  $\epsilon^{mat}$  and inclusions with permittivity  $\epsilon^{inc}$  (e.g Wolfensberger and Berne (2018)):

$$\epsilon_{eff}(\epsilon^{mat}, \epsilon^{inc}, f_{vol}^{inc}) = \epsilon^{mat} \left( \frac{1 + 2f_{vol}^{inc} \frac{\epsilon^{inc} - \epsilon^{mat}}{\epsilon^{inc} + 2\epsilon^{mat}}}{1 - f_{vol}^{inc} \frac{\epsilon^{inc} - \epsilon^{mat}}{\epsilon^{inc} + 2\epsilon^{mat}}} \right)$$
(31)

with  $f_{vol}^{inc}$  the volume fraction of the inclusions.

• Dry snow (DIEL=3), dry graupel (DIEL=7) and primary ice (DIEL=11) Dry solid hydrometeors consist of inclusions of ice in a matrix of air. In this case  $\epsilon^{mat} \approx 1$ , and  $f_{vol}^{inc} = \frac{V_i}{V_d} = \frac{\rho_d}{\rho_i}$ .

The permittivity of the dry hydrometeor  $\epsilon_d$  can be estimated following equation (4) in Ryzhkov *et al.* (2011)):

$$\epsilon_d = \frac{1 + 2\frac{\rho_d}{\rho_i}\frac{\epsilon_i - 1}{\epsilon_i + 2}}{1 - \frac{\rho_d}{\rho_i}\frac{\epsilon_i - 1}{\epsilon_i + 2}}$$
(32)

where  $\rho_d$  is the density of the dry hydrometeor (snow, graupel or ice crystals),  $\epsilon_i$  the dielectric constant of pure ice (estimated from Hufford (1991)), and  $\rho_i$  the density of pure ice (0.916 g cm<sup>-3</sup>).

• Wet graupel

## DIEL=7

The melting process of the graupel implemented is summarized in Figure 2. In its initial state, the particle is dry and of fairly low density (between 100 and 200 kg m<sup>-3</sup> depending on its size). When melting, the water first fills the outer cavities of the particle, forming a layer which is treated as an ice matrix with water inclusions. This outer layer plays the role of matrix for the whole particle while the core (still an air matrix with ice inclusions) plays the role of an inclusion. When the graupel is fully soaked with water ( $F_w = F_{w_{sat}}$ ) it is considered as an ice matrix with water inclusions. The latter is then represented as a matrix of liquid water with a water-saturated graupel inclusion.

In summary, when the water fraction increases, the particle can be successively considered as a matrix of water-saturated graupel (matrix of ice with water inclusions) with dry graupel inclusions (matrix of air with ice inclusions) and as a matrix of water with water-saturated graupel (matrix of ice with water inclusions). Then, we can express the effective permittivity as follows:

$$\epsilon_{wg} = \begin{cases} \epsilon_{eff}(\epsilon_{eff}(\epsilon_i, \epsilon_w, f_{vol}^{iw}), \epsilon_{eff}(\epsilon_a, \epsilon_i, f_{vol}^{ai}), f_{vol}^{iw} - a^{ai}) &, \text{ if } F_w \le F_{w_{sat}} \\ \epsilon_{eff}(\epsilon_w, \epsilon_{eff}(\epsilon_i, \epsilon_w, f_{vol}^{iw}), f_{vol}^{w} - a^{iw}) &, \text{ if } F_w > F_{w_{sat}} \end{cases},$$
(33)

where  $f_{vol}^{iw}$ ,  $f_{vol}^{ai}$ ,  $f_{vol}^{iw}$  and  $f_{vol}^{w}$  are the volume fractions of the corresponding inclusions.

 $f_{vol}^{iw}$  is the volume fraction of the cavities. It is constant all along the early stage of the melting process ( $F_w \leq F_{w\_sat}$ ). Consequently,  $f_{vol}^{iw}$  is equal to the ratio between the initial volume of the cavities (see Equation 15) and the initial volume of the particle:

$$f_{vol}^{iw} = \frac{V_c^{(i)}}{V_g} = 1 - \frac{\rho_g}{\rho_i}$$
(34)

 $f^{ai}_{vol}$  is the volume fraction of ice in the initial particle:

$$f_{vol}^{ai} = \frac{V_i}{V_g} = \frac{\rho_g}{\rho_i} \tag{35}$$

The cavities are supposed to be evenly distributed within the graupel. Thus, the volume fraction of the core (made of ice and air) within the particle  $f_{vol}^{iw}$  and be expressed as the ratio between the volume of the unfilled cavities in the core  $V_c - V_w$  and the total volume of the cavities  $V_c$ :

$$f_{vol}^{iw}{}^{ai} = \frac{V_c - V_w}{V_c} \tag{36}$$

Combining equations 17, 18 and 19, we can retrieve:

$$f_{vol}^{iw} = 1 - \frac{D^3}{D_{eq}^3} \frac{F_w \rho_g}{\rho_w \left(1 - \frac{\rho_g}{\rho_i}\right)}$$
$$f_{vol}^{iw} = 1 - \frac{F_w}{1 - F_w} \frac{\rho_i \rho_g}{\rho_i \rho_w - \rho_g \rho_w}$$
(37)

 $f_{vol}^{w\_iw}$  is the volume fraction of the inner core of the particle (ice + water filled cavities) once  $F_w > F_{w\_sat}$ .

The volume of the core  $V_{core}$  can be estimated by considering the mass of ice within the core, which is equal to the total mass of ice of the particle:

$$m_i = (1 - F_w)m_g$$

 $m_i$  can be also expressed as a function of the volume of the core  $V_{core}$  if we consider that the water is removed from the cavities:

$$m_i = \rho_g V_{core}$$
$$V_{core} = (1 - F_w) \frac{m_g}{\rho_g}$$
$$= (1 - F_w) V_q$$

$$V_{core} = (1 - F_w)\frac{\pi}{6}D^3$$
(38)

$$f_{vol}^{w\_iw} = \frac{V_{core}}{V_{wg}}$$
$$= \frac{(1 - F_w)\frac{\pi}{6}D^3}{\frac{\pi}{6}D_{eq}^3}$$

$$f_{vol}^{w\_iw} = \frac{(1-F_w)}{\rho_g \left(\frac{1-F_w}{\rho_i} + \frac{F_w}{\rho_w}\right)}$$
(39)

In summary, the volume fractions involved in the calculation of the effective permittivity of the melting graupel are:

$$\begin{aligned}
f_{vol}^{iw} &= 1 - \frac{\rho_g}{\rho_i} \\
f_{vol}^{ai} &= \frac{\rho_g}{\rho_i} \\
f_{vol}^{iw} - a^{ai} &= 1 - \frac{F_w}{1 - F_w} \frac{\rho_i \rho_g}{\rho_i \rho_w - \rho_g \rho_w} \\
f_{vol}^w &= \frac{(1 - F_w)}{\rho_g \left(\frac{1 - F_w}{\rho_i} + \frac{F_w}{\rho_w}\right)}
\end{aligned} \tag{40}$$

## 2.2.6 Compiling

• Before compiling:

for calculations for one hydrometeor type only, modify the loop over types:

DO idtype=1,1 for rain only, for example.

for calculations for one radar frequency band only, modify the loop over bandes:

DO idbande=3,3 for S-band only, for example.

• Compile with make: builds the executable Tmat

#### 2.2.7 Running the program

- make sure an empty directory OUTPUT is present in DPOLSIMUL (otherwise create this directory: mkdir OUTPUT)
- make sure that the file ampl.par.f is present (in the same directory)
- make sure that the parameter files corresponding to the frequency band defined (band="S" for example) are present in the same directory. For this example, files: TmatParam\_Swg, TmatParam\_Sgg, TmatParam\_Srr, TmatParam\_Sss, TmatParam\_Sii
- Launch the command ./Tmat : creation of the following files for each hydrometeor type (example for rain here) :
  - TmatVarPol Srr : contains
    - \* 3 first lines describing the content of the file : LAMmin LAMstep LAMmax ELEVmin ELEVstep ELEVmax Tcmin Tcstep Tcmax expDmin expDstep expDmax Fwmin Fwstep Fwmax
    - \* the dualpol variables for a range of wavelengths, temperatures, elevation angles, liquid water fractions (Fw) in case of type w (melting graupel), and diameters: LAM, Tc, ELEV, Fw, D, Deq, ZhhlgR, ZvvlgR, Zhhlg, Zvvlg, Zdrlg, Rhv, KDP, KdpR, Adp, Ah, Av, Deltaz
  - TmatCoefDiff\_Srr : contains

- \* 3 first lines describing the content of the file : LAMmin LAMstep LAMmax ELEVmin ELEVstep ELEVmax Tcmin Tcstep Tcmax expDmin expDstep expDmax Fwmin Fwstep Fwmax
- \* the real and imaginary parts of the scattering coefficients (f for forward scattering), and Zhh, Zvv et Kdp calculated with Rayleigh for spheroid scattering method:

LAM, Tc, ELEV, Fw, D, Dm, Deq, Deqm, Deqr, Deqrm, REAL(S11), AIMAG(S11), REAL(S22), AIMAG(S22), REAL(S11f), AIMAG(S11f), REAL(S22f), AIMAG(S22f), ZhhR, ZvvR, KdpR

 TmatResu\_Srr : this file contains just some information that the user can select to print if needed (like the dielectric constant)

## 3 Program TmatInt.f90

Reading of scattering coefficients files (like TmatCoefDiff\_Srr) produced by the previous program and integration over diameters (from the particle size distribution PSD) in order to build new tables with integrated coefficients.

#### 3.1 Steps of the program

- 1. Loops over the band and the hydrometeor type
- 2. Reading of the hydrometeor constants file for PSD calculations + the rain constant file (for the equivalent melted hydrometeor)
- 3. Reading of the min/step/max for LAM, ELEV, Tc, expD and Fw in the 2nd line of the scattering coefficients file
- 4. Reading of the coefficients RES11, IMS11, RES22, IMS22, RES11f, IMS11f, RES22f, IMS22f for each values of LAM, ELEV, Tc, expD and Fw
- 5. Loops over wavelengths LAM, temperature Tc, elevation ELEV, 3d parameter of the table P3 (liquid water fraction Fw if 1-moment, concentration CC if 2-moments), and hydrometeor content M
  - (0) Step 0: calculation of the "liquid" and "solid" parts of the contents  $(M_{rr}$  and  $M_{ss})$  with a Riemman sum (integration over diameters)

The purpose of this step is to calculate a first estimation of the "liquid" and "solid" parts of the distribution  $(N_{rr} \text{ and } N_{ss})$  and of the "liquid" and "solid" parts of the hydrometeor content when doing the Riemann sum  $(M_{rr} \text{ and } M_{ss})$ .

- loop over diameters (D in mm)
  - calculation of the position (kTmat) of the parameters (LAM, Tc, ELEV, P3, M, D) in the table
  - reading of all diameters and coefficients in the table
  - number concentration of the "solid part"  $N_{ss} = PSD_{ss}[(1 F_w)M, D_m],$ and "liquid part" of the distribution  $N_{rr} = PSD_{rr}(F_wM, D_{eqr_m})$

- Liquid  $M_{rr}$  and solid  $M_{ss}$  parts of the content (in case of a wet specie):

$$M_{rr} = \sum_{D_{min}}^{D_{max}} a_{rr} D_{eqr_m}^{b_{rr}} N_{rr} (D_{eqr_m}) \times 0.5 (D_{eqr_{msup}} - D_{eqr_{minf}})$$
(41)

$$M_{ss} = \sum_{D_{min}}^{D_{max}} a_{ss} D_m^{b_{ss}} N_{ss}(D_m) \times 0.5 (D_{m_{sup}} - D_{m_{inf}})$$
(42)

#### (1) Step 1: correction of $N_{rr}$ and $N_{ss}$

To ensure the correspondence between the initial content M and the recalculated content  $M_{int}$  after integration with the Riemann sum, a corrective factor  $\frac{M}{M_{int}}$  is applied to the particle numbers  $N_{rr}$  and  $N_{ss}$  before computation of the radar variables through integration over the PSD.

- loop over diameters (D in mm)
  - calculation of the position (kTmat) of the parameters (LAM, Tc, ELEV, P3, M, D) in the table
  - reading of all diameters and coefficients in the table
  - estimation of the fall velocity of the wet specie  $v_w$  by combining the rain and solid hydrometeor fall velocities ( $v_{rr}$  and  $v_{ss}$ ): eq 4.15 p 93 in Le Bastard (2019), following Wolfensberger and Berne (2018) and Mitra (1990)
  - Application of the corrective factor to the number concentrations:

\* 
$$N_{ss}$$
 is multiplied by  $\frac{(1 - F_w)M}{M_{ss}}$   
\*  $N_{rr}$  is multiplied by  $\frac{F_wM}{M_{rr}}$ 

- Combination of the liquid and solid parts of the distribution following Szyrmer and Zawadzki (1999) and used by Wolfensberger and Berne (2018) and Le Bastard (2019) (see their equation 4.20 p 107) to compute the total particle number concentration  $N(D_{eq})$ :

$$N(D_{eq}) = (1 - F_w) \frac{v_{ss}}{v_w} N_{ss}(D_{eq}) + F_w \frac{v_r}{v_w} N_{rr}(D_{eq})$$
(43)

Another formulation for the estimation of the number concentration for wet hydrometeors (called the "weighted PSD approach") is also proposed by Wolfensberger and Berne (2018) with their equation (41). This formulation has not yet been implemented in the T-matrix code.

- Recalculation of the total content  $M_{int}$  after integration using the total particle number concentration  $N(D_{eq})$ 

$$M_{int} = \sum_{D_{min}}^{D_{max}} a_{rr} D_{eqr_m}^{b_{rr}} N(D_{eq}) \times 0.5 \left( D_{eq_{m_{sup}}} - D_{eq_{m_{inf}}} \right)$$
(44)

#### (2) Step 2: correction of the total particle number concentration N

- loop over diameters
  - re-computation of  $N_{rr}$  and  $N_{ss}$  and application of the corrective factors as in Step 1
  - re-computation of  $N(D_{eq})$  as in Step 1

- application of the corrective factor  $\frac{M}{M_{int}}$  to  $N(D_{eq})$
- computation of the integrated rainfall rate  $RR_{int}$
- integration of the scattering coefficients over the PSD:
- integration of Zhh, Zvv et Kdp calculated with Rayleigh for spheroid scattering
- writing of the integrated coefficients in the output file: TmatCoefInt\_ICE3\_S106.2\_rr (example for ICE3, S band and rain)
- computation of the radar variables (but for a single hydrometeor type)  $Z_{hh}$ ,  $Z_{dr}$ ,  $K_{dp}$ ,  $\rho_{hv}$ ,  $A_{dp}$  (differential attenuation),  $A_h$  and  $A_v$  (specific attenuation on the H and V polarizations),  $\delta_{hv}$  (back-scattering differential phase)
- writing of the radar variables in the output file (for example TmatVarInt\_ICE3\_S106.2\_rr)

#### 3.2 Compiling

- before compiling, choose the frequency band in the beginning of the file (they should be the same as those selected in the previous program)
- compiling with the command make -f makeTmatInt => build the executable TmatInt

#### 3.3 Running the program

- INPUT: you need to have files like TmatCoefDiff\_Sgg, for each hydrometeor type
- RUN: ./TmatInt
- OUTPUT: files TmatCoefInt\_Srr and TmatVarInt\_Srr with integrated scattering coefficients and dualpol variables
  - TmatCoefInt\_Srr: Tc, ELEV, P3, M, S11carre, S22carre, REAL(S22S11), AIMAG(S22S11), ReS22fmS11f, ImS22ft, ImS11ft, RRint
  - TmatVarInt\_Srr:
     Tc, ELEV, P3, M, Zhhlg, Zvvlg, Zdrlg, Rhv, KDP

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